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1 **A numerical study of capillary pressure-saturation relationship for supercritical**
2 **carbon dioxide (CO₂) injection in deep saline aquifer**

3
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8
9 **Abstract**

10 Carbon capture and sequestration (CCS) is expected to play a major role in reducing
11 greenhouse gas in the atmosphere. It is applied using different methods including geological,
12 oceanic and mineral sequestration. Geological sequestration refers to storing of CO₂ in
13 underground geological formations including deep saline aquifers (DSAs). This process
14 induces multiphase fluid flow and solute transport behaviour besides some geochemical
15 reactions between the fluids and minerals in the geological formation. In this work, a series of
16 numerical simulations are carried out to investigate the injection and transport behaviour of
17 supercritical CO₂ in DSAs as a two-phase flow in porous media in addition to studying the
18 influence of different parameters such as time scale, temperature, pressure, permeability and
19 geochemical condition on the supercritical CO₂ injection in underground domains. In contrast to
20 most works which are focussed on determining mass fraction of CO₂, this paper focuses on
21 determining CO₂ gas saturation (i.e., volume fraction) at various time scales, temperatures and
22 pressure conditions taking into consideration the effects of porosity/permeability, heterogeneity
23 and capillarity for CO₂-water system. A series of numerical simulations is carried out to illustrate
24 how the saturation, capillary pressure and the amount of dissolved CO₂ change with the change
25 of injection process, hydrostatic pressure and geothermal gradient. For example, the obtained
26 results are used to correlate how increase in the mean permeability of the geological formation
27 allows greater injectivity and mobility of CO₂ which should lead to increase in CO₂ dissolution
28 into the resident brine in the subsurface.

29
30 **Keywords:** geological sequestration, two-phase flow, capillary pressure, porous media, CO₂
31 sequestration, deep saline aquifer, CO₂ sequestration

32
33 **1. Introduction**

34 Carbon sequestration is a technique for managing carbon dioxide (CO₂) that has been emitted
35 into the atmosphere by various activities, e.g., combustion of carbon-based fuels. It is a
36 relatively new concept that had been developed to address the problem of global warming,

37 which is attributed to high levels of atmospheric CO₂. In a more specific approach, geological
38 sequestration aims to inject supercritical CO₂ into porous formations underground while
39 attempting to prevent leakage of CO₂ to the surface again. This method can be applied to
40 declining oil fields, un-minable coal seams as well as deep saline aquifers (DSAs). Injecting
41 CO₂ into DSAs is considered to be one of the most feasible sequestration methods of CO₂.
42 From a fluid mechanics point of view, injecting supercritical CO₂ into geological formations can
43 be treated as a two-phase flow in a porous medium (Tsang et al., 2008). Supercritical CO₂ is
44 considerably denser than the gaseous CO₂ phase but has lower density and viscosity than the
45 occupant brine in the porous space. As a result of the differences of fluid densities,
46 supercritical CO₂ migrates buoyantly towards the upper confining layer. The preferred depths
47 to inject CO₂ are greater than 800m (Prevost et al., 2005) as they provide the required
48 conditions above the critical points of CO₂ for it to stay in supercritical phase. This increases
49 the storage capacity of the site because more CO₂ can be stored within a specific volume.

50

51 It must be emphasised that there are particular conditions, which the geological formation must
52 meet for CO₂ storage to be successful. According to Bachu and Bennion (2008), three basic
53 conditions must be met, namely, (i) capacity, i.e., the geological media must have the capacity
54 to allow the anticipated amount of CO₂ over the duration of the project operation; (ii) injectivity,
55 i.e., the media must be able to allow the CO₂ at its injection rate and, (iii) confinement, i.e., the
56 media must be able to impede leakage of CO₂ from the storage zone or minimize leakage to the
57 tolerable levels. Furthermore, geological storage of CO₂ is determined by four foremost trapping
58 mechanisms as discussed below.

59 (a) Structural trapping, which takes place when CO₂ gas becomes immobile in the porous
60 sedimentary layers with existed brine by impermeable barriers (White et al., 2013).

61 (b) Residual trapping that takes place as a result of the hysteresis effect when the
62 saturation direction is reversed after the injection process stops and, the existing brine
63 moves back and tries to displace CO₂ in the pores (Ide et al., 2007).

64 (c) Solubility trapping takes place when the injected CO₂ dissolves in the resident fluid and
65 increases the acidity and density of the brine creating convective currents that allow the
66 denser brine with high concentration of CO₂ to settle at the bottom part of the aquifer
67 trapping the CO₂ more securely (Silin et al. 2009).

68 (d) Mineral trapping occurs when the dissolved CO₂ reacts with the brine producing
69 carbonic acid that reacts with the dissolved ions within the aquifer brine and minerals
70 forming the host rock resulting in chemical precipitation of solid carbonate minerals
71 (Beni et al., 2012).

72

73 Modelling of underground injection of CO₂ primarily represents modelling a system of two-
74 phase flow in porous media which requires one to identify the relevant parameters. These
75 parameters describe various physical and chemical properties of the geological formation such
76 as entry pressure (depending on pore/particle size of the domain), hydrodynamic conditions
77 (e.g., pressure difference, groundwater velocity), fluid properties, permeability, chemical species
78 from geochemical reactions and fluid/fluid interfacial mass transfer (Ide et al., 2007).
79 Considerable uncertainty may however exist with regards to the formation-related parameters
80 because of the difficulty in collecting sufficient data across the huge areas that should be taken
81 into account for any geologic sequestration project. A number of studies have been conducted
82 to determine the capillary pressure-saturation-relative permeability relationships for subsurface
83 injection of CO₂ into porous media (e.g., Bachu et al., 1994; Pruess et al., 2003; Kumar et al.,
84 2005; Knauss et al., 2005; Juanes et al., 2006; Birkholzer et al., 2009; Schnaar and Digiulio,
85 2009). They demonstrate that computational models are able to replicate complex formation
86 heterogeneities by employing statistical routines; residual CO₂ trapping and hysteretic relative
87 permeability curves, dissolution reactions and mineral precipitation and others. For example,
88 Nordbotten et al. (2004) analytically described the time evolution of the CO₂ plume dominated
89 by viscous forces with irrelevant effects of the CO₂ buoyancy forces using a simplified form of
90 Buckley-Leverett equation. They utilized their modelling results to inspect the
91 accuracy/implication of assuming constant properties for the fluids in the storage formation and
92 discussed some cases where buoyancy and non-zero residual saturations have more influence
93 on the mobility of CO₂ plume in addition to the effects of CO₂ dissolution in the existing brine.

94
95 One of the critical issues in CO₂ geological sequestration is the phase transition from liquid or
96 supercritical to gas according to the temperature and pressure changes during the injection
97 progression. Therefore, numerical simulation for CO₂ sequestration in saline aquifers should
98 have the ability to envisage when the CO₂ phase transition occurs. It must also be able to
99 determine the buoyancy and viscous forces influence on the fluid flow and, CO₂ dissolution in
100 the aqueous fluid (White and Oostrom, 2003). Though capillarity plays a crucial role in
101 assessment of saline aquifers for CO₂ sequestration, there is not much real (field) data
102 available about the behaviour of CO₂-brine flow system in the porous rocks. Plug and Bruining
103 (2007) developed a laboratory scale method to investigate the static capillary pressure change
104 as a function of saturation at different pressure and temperature conditions. They examined the
105 influence of CO₂ dissolution in water by comparing its behaviour to the behaviour of nitrogen (N₂)
106 under the same conditions and observed that the residual water saturation (S_{wc}) for CO₂ is
107 much smaller than that for N₂ due to the difference in interfacial tension.

108

109 Capillary trapping, which is also called residual trapping, is closely related to the capillary forces
110 between CO₂ and resident fluid at the scale of the grains of reservoir rock which is controlled by
111 interfacial forces, pore size and wettability (Alkan et al., 2010). Experimental studies conducted
112 by Bennion and Bachu (2008), and Plug and Bruning (2007) reported that permeability and
113 capillarity are influenced by interfacial forces and wettability of CO₂-brine-rock systems. Another
114 study on CO₂ sequestration has been conducted by Bickle et al. (2007) who modelled CO₂ flow
115 behaviour in Sleipner field in the North Sea. They used a theoretical model and validated it with
116 experimental results by Lyle et al. (2005) to characterise the gravity flow in porous media. To
117 attain their solutions they employed a number of assumptions, e.g., neglecting the motion of the
118 existing fluid within the hosting formation and ignored both capillary and viscous forces in the
119 fluid flow system which exhibited some limitations in the applicability of their solutions. Bickle et
120 al. (2007) concluded that the radius of accumulated CO₂ ponds in the subsurface increases
121 linearly with the square-root of the elapsed time. They observed an increase in CO₂ input in
122 higher layers of the domain with a decrease in lower ones due to the leakage into the upper
123 structures of the modelled formation. Their solutions provide important predictions on CO₂
124 behaviour with no need to carry out full simulation for any potential storage sites.

125
126 Unlike most conventional approaches for determining capillary pressure relationships which are
127 based on equilibrium flow conditions (i.e., time derivative of fluid saturation is zero), dynamic
128 capillary pressure effects have been shown to have a great influence on two-phase flow in
129 porous media (Helmig et al., 2007; Mirzaei and Das, 2007; Hanspal and Das, 2012). A number
130 of fundamental studies (e.g., Oung et al., 2005; Manthey et al., 2005; Bottero et al., 2006) have
131 investigated the dynamic capillary pressure effects in two-phase flow systems, and this gives
132 rise to the possibility of applying these understanding to determine if these effects are
133 significant for supercritical CO₂ flow in the geological formation as well.

134
135 In addressing most of the above issues, the main goal of this study is to carry out a simulation
136 study to determine static and dynamic capillary pressures for CO₂-water system as a function of
137 saturation for different permeability and heterogeneity at various time scales, temperature and
138 pressure conditions in order to evaluate the implications of different CO₂ injection strategy and
139 its storage capacity in briny aquifers. For this purpose a series of numerical simulations are to
140 be carried out under various pressure, temperature, heterogeneity and injection rate conditions.
141 It is envisaged that this would help the prediction of the right CO₂ injection process and CO₂
142 behaviour within the aquifer formation during sequestration life time which has a vast impact on
143 the energy cost and storage process safety. It is believed that this study will provide better
144 understanding of the injection and sequestration processes.

145

146 2. Modelling Approach

147

148 2.1. Main Equations

149 In this work the injection of CO₂ into saline aquifers is defined to represent the flow of two
150 immiscible fluids, namely, water (brine) as a wetting phase and CO₂ as a non-wetting phase in
151 a porous medium where supercritical CO₂ replaces the existing fluid in a process called
152 drainage.

153

154 2.1.1. Mass and Momentum Conservation Equations

155 Modelling CO₂ injection into geological formation is governed by the equations of mass and
156 momentum conservation.

157

158 The conservation of momentum is described by the following form of Darcy's law:

$$159 \frac{\partial(S_\alpha \phi \rho_\alpha)}{\partial t} + \nabla \cdot (\rho_\alpha v_\alpha) - \rho_\alpha q_\alpha = 0 \quad (1)$$

160 where S_α is the phase α (water or CO₂) saturation, ϕ is the porosity, ρ_α is the density, t refers to
161 the elapsed time, v_α is the average pore velocity of the phase and q_α refers to the phase flux.

162

163 From the generalized Darcy's law (equation of momentum), velocity vector v_α is calculated as

$$164 v_\alpha = - \frac{k_{r\alpha}}{\mu_\alpha} K (\nabla p_\alpha - \rho_\alpha g) \quad (2)$$

165 $k_{r\alpha}$ identifies the relative permeability for the phase α (water or CO₂), μ refers to the dynamic
166 velocity, p_α identifies the pressure, K is the tensor of absolute permeability (defined to be
167 isotropic) and g the vector of gravity. The phase permeability (effective permeability) (k_α) is
168 related to the relative permeability ($k_{r\alpha}$) as:

$$169 k_{r\alpha} = \frac{k_\alpha}{K} \quad (3)$$

170 where K signifies the domain permeability for a single-phase flow (the absolute permeability).

171

172 By substituting equation (2) in equation (1) the following general form of mass conservation
173 equation is obtained for both fluid phases:

$$174 \frac{\partial(S_\alpha \phi \rho_\alpha)}{\partial t} - \nabla \cdot \left(\rho_\alpha \frac{k_{r\alpha}}{\mu_\alpha} K (\nabla p_\alpha - \rho_\alpha g) \right) - \rho_\alpha q_\alpha = 0 \quad (4)$$

175

176 2.1.2. Constitutive Relationships

177 The two fluid flow process is dominated by capillary pressure (P_C) - saturation (S_w) - relative
178 permeability (k_r) relationships because any decrease in the wetting phase saturation results in
179 non-wetting fluid retreatment into smaller pores which increases the capillary pressure. In a two

180 phase flow the capillary pressure is defined as the difference between the average phase
 181 pressures of non-wetting (nw) and wetting (w) phases,

$$182 \quad P_c = P_{nw} - P_w \quad (5)$$

183

184 One of the most common formulations used to determine P_c - S_w - K_r relationships is Brooks-
 185 Corey function (Brooks and Corey, 1964), in which the displacement pressure of the wetting
 186 fluid from the largest pore (P_d) is involved while this pressure has been ignored by other authors
 187 for fully saturated porous media. The relationship defines the effective saturation as:

$$188 \quad S_{ew} = \left(\frac{P_c}{P_d} \right)^{-\lambda} \quad \text{for } P_c > P_d \quad (6)$$

$$189 \quad S_{ew} = 1 \quad \text{for } P_c \leq P_d \quad (7)$$

$$190 \quad S_{ew} = \frac{(S_w - S_{wr})}{(1 - S_{wr})} \quad \text{for } 0 \leq S_{ew} \leq 1 \quad (8)$$

191 Where, (S_{ew}) denotes the effective water saturation, (S_{wr}) is the residual water saturation, (P_d)
 192 represents entry (displacement) pressure, (λ) is the pore size distribution index.

193

194 Brooks-Corey correlations in conjunction with the Burdine theorem (Burdine, 1953) are used to
 195 define the relative permeability-saturation relationships for wetting (w) and non-wetting (nw)
 196 phases.

$$197 \quad k_{rw} = S_{ew}^{\frac{2+3\lambda}{\lambda}} \quad (9)$$

$$198 \quad k_{rnw} = (1 - S_{ew})^2 (1 - S_{ew}^{\frac{2+\lambda}{\lambda}}) \quad (10)$$

199

200 The coupled equations are solved for the primary variables where the porous domain is
 201 assumed to be a rigid rock and both fluids are defined as incompressible. Furthermore, the
 202 dynamic viscosities of the fluids are assumed to be constant and all source and sink terms are
 203 ignored.

204

205 **2.2. Simulation Approach**

206 The scope of this research is to simulate the process of injecting CO_2 as a supercritical fluid into
 207 DSAs. It focuses on the flow of multiphase fluid (H_2O - CO_2 - NaCl) in a porous media for which
 208 STOM32 (STOMP- CO_2) operational mode of STOMP (*subsurface transport over multiple*
 209 *phases*) simulation code is used. In this mode water (brine) is the wetting phase and CO_2 is a
 210 non-wetting fluid which is injected at different pressure rates into the porous domain which is
 211 fully saturated with water (brine). This leads to a situation where CO_2 drains water out of the
 212 domain in a process called drainage followed by an imbibition process when CO_2 injection ends

213 and water flows back into the domain to replace CO₂ in the domain pores leaving some traces
214 of it trapped. This operational mode is able to incorporate buoyancy and viscous forces driven
215 flow, CO₂ dissolution in aqueous fluid, phase transition, dispersion and diffusivity of the gas and
216 uses the finite volume technique to numerically simulate the process. These are discussed in
217 detail by White and Oostrom (2003) and are not repeated in this paper. However it should be
218 mentioned that STOMP-CO₂ simulator is written in FORTRAN 90 with a capability of dynamic
219 memory allocation for faster execution. The collection of source files is required to be compiled
220 into an executable file that can be used on various computing platforms including Linux and
221 Windows to read the input file that is created by the user including a number of cards that
222 contain calculation instructions and required parameters to solve the simulation problem. The
223 code has been effectively optimized for workstations (HP, IBM, Sun) in addition to mainframes.
224 The speed and memory requirements for running STOMP-CO₂ executable files depend on the
225 complexity of the problem and computational grid refinement. There is no minimum memory or
226 processor speed provided by the developer. However, from our experience it has been found
227 out that the code better functions on UNIX operating system with 2.4 GHz CPU and 1 GB
228 memory. STOMP-CO₂ is utilized to numerically solve the coupled conservation equations (water
229 mass, CO₂ mass and NaCl mass) by converting them to algebraic equations using finite volume
230 method (FVM) and Euler-backward time differencing for spatial and temporal discretizations,
231 respectively. Backword Euler method is a first order time stepping method that makes an error
232 of Δt^2 for each time step. This method offers more stability and accuracy than forward Euler
233 method especially for problems with large and nonlinear functions like diffusion equations. The
234 produced algebraic equations in the discretised equations are closed using a number of
235 constitutive relationships as explained in section 2 and solved using Newton-Raphson iteration
236 to resolve their nonlinearities (White and Oostrom, 2003).

237

238 **2.3. Initial and Boundary Conditions**

239 The domain is considered to be anisotropic and almost fully saturated with brine before injecting
240 supercritical CO₂ in the centre. The initial condition for all simulation conditions are shown in
241 Table 4. We generate two-phase conditions within the computational domain by setting the
242 aqueous saturation value at 0.9999 as an initial condition for the employed equations of state in
243 the simulation code (e.g., Kelvin equation (Nitao, 1988), and the formulation by Battistelli et al.
244 (1997)) which take into account the changes in thermodynamic properties of the fluid phases as
245 the simulation conditions change. The non-wetting fluid (CO₂) saturation was assumed to be
246 1.0 at the injection source at the outer wall of the reservoir and 0.00001 in the rest of the
247 computation domain as initial condition for the reason above. It is injected into the lower 3 grid
248 cells (i.e. 30 m from the bottom of the domain). Vertically zero flux is considered for aqueous
249 phase at the well case as inner boundary while the outer boundary was assumed to be infinite

250 with zero flux for gas phase. Horizontally zero flux is considered at the upper and lower
 251 surfaces which force the injected gas to spread laterally. For both dynamic and static conditions,
 252 fluids saturation, pressure and volume are measured at each node and the CO₂ saturation is
 253 plotted versus simulation time. This procedure is repeated twice: once for sandstone (coarse)
 254 and another for Wechselfolgen (fine) homogeneous domains. This procedure is repeated twice:
 255 once for sandstone (coarse) and another for Wechselfolgen (fine) homogeneous domains.

256

257 **2.4. Dynamic and Quasi-static Simulations**

258 In this research work, simulations are carried out by injecting CO₂ into the centre of the
 259 computational domain which is initially fully saturated with brine. The gas pressure is defined to
 260 be zero all over the domain. The CO₂ injection starts at 32 MPa and increased at a rate of 0.1
 261 MPa every 0.5 year for 20 years for quasi-static simulations. This increment in injection
 262 pressure increases the capillary pressure (P_c) in the domain until it reaches the displacement
 263 pressure (P_d) when the injected CO₂ starts displacing the existing brine and continues till a
 264 steady state is reached when average values of aqueous saturation and capillary pressure are
 265 calculated to give a single point for the $P_c - S_w$ relationships. This procedure is repeated for
 266 each time step from which the $P^c - S_w$ curves are produced. For dynamic simulations the
 267 imposed injection pressure is increased to 36 MPa in one step and maintained till the end of
 268 injection period.

269

270 **2.5. Capillary Pressure and Saturation Averaging**

271 From the locally predicted values of saturation and pressure at each grid node for each time
 272 step (t_n) the volume-weighted average water saturation (S_w) and saturation-weighted average
 273 capillary pressure (P_c) values for the whole domain are determined using the following
 274 equations.

275

276 The average saturation at any time step (t_n) is calculated by;

$$277 \quad S_w |_{t_n} = \frac{\sum_{j=1}^m S_{wj} V_j |_{t_n}}{\sum_{j=1}^m V_j} \quad (11)$$

278

279 And the average capillary pressure is calculated by;

$$280 \quad P^c |_{t_n} = \left[\frac{\sum_{j=1}^m (1-S_{wj}) P_{nwj}}{\sum_{j=1}^m (1-S_{wj})} - \frac{\sum_{j=1}^m S_{wj} P_{wj}}{\sum_{j=1}^m S_{wj}} \right] |_{t_n} \quad (12)$$

281

282 Where, V_j , is the volume of node j , and, S_{wj} , P_{wj} and P_{nwj} denote water saturation, water
 283 pressure and CO₂ pressure at node j , respectively.

284

285 The time derivative of saturation dependency can be calculated from the average saturation
 286 values calculated from equation (13) as follows:

$$287 \quad \left. \frac{\partial s}{\partial t} \right|_{s_w t_n} = \frac{S_w|_{t_{n+1}} - S_w|_{t_{n-1}}}{t_{n+1} - t_{n-1}} \quad (13)$$

288
 289 As shown in equations (11) and (12), both calculated average values are based on water
 290 saturation and, hence, they are called saturation-weighted averages (Mirzaei and Das, 2007;
 291 Hanspal and Das, 2012).

292
 293 Conventional theories (Collins, 1961; Scheidegger, 1974; Bear and Verruijt, 1987; Helmig, 1997)
 294 define capillary pressure as a function of fluid saturation only for fluids at equilibrium conditions.
 295 However, this is not always the case as fluids might not flow under steady conditions especially
 296 at early stages of flow when the change rate of saturation is thought to be high. Therefore, it
 297 has been suggested by many authors that an additional term ought to be added to the capillary
 298 pressure equation (5) for dynamic fluid flow in porous media (Hassanizadeh and Gray, 1993a;
 299 Beliaev and Schotting, 2002; Dahle et al., 2005; Hanyga and Serebnynska, 2005; Oung et al.,
 300 2005).

301
 302 In this study we will be investigating the dynamic effects at a field-scale domain. The additional
 303 term is called the dynamic coefficient (τ) which represents dynamic capillary pressure effect on
 304 the flow behaviour and is determined from the slope of a linear relationship between the
 305 capillary pressures at dynamic and static flow conditions and the time derivative of saturation as
 306 shown in equation (14):

$$307 \quad (P_{c,dyn} - P_{c,stat})|_s = -\tau \left. \frac{\partial s}{\partial t} \right|_s \quad (14)$$

308 where $P_{c,dyn}$ and $P_{c,stat}$ represent dynamic and static capillary pressures calculated at a specific
 309 value of saturation (s), respectively. The dynamic coefficient has been used by many previous
 310 workers (e.g. Tian et al., 2012; Fucik, 2010; Mirzaei and Das, 2007; Hanspal and Das, 2012;
 311 Das and Mirzaei, 2013; Mirzaei and Das, 2013; Hanspal et al., 2013) to take into account
 312 dynamic capillary pressure effect and, therefore, a detailed discussion on dynamic capillary
 313 pressure effect is avoided in this paper.

314
 315
 316 **2.6. Computational Domain**

317 Our simulation parameters are based on Bunter Sandstone Aquifer in North German Basin in
 318 North-Eastern Germany. This aquifer consists of four cycles beginning with basal sandstone
 319 which has three cycles of permeable layers (Detfurth, Hardegsen and Solling-Folge) and ending

320 with an alternating succession of silt, sand and clay stone (May et al., 2004). In this research
321 study we focus on Detfurth cycle which is divided to a lower sandstone with high permeability
322 and upper alternating succession of sand, silt and clay stones which is called (Wechselfolgen)
323 with low permeability because it demonstrates heterogeneity in regards to porosity and
324 permeability.

325
326 The simulated three-dimensional cylindrical domain extends laterally (r-direction) from the
327 injection point which is represented by the well radius of 0.2 to 2500 m and vertically from 2900
328 to 3000 m below land surface while at the top and bottom are impermeable layers that preserve
329 the injected CO₂ safely in the storage formation. This depth ensures that the injected CO₂ will
330 remain in supercritical state which increases the storage capacity of the site. The system can be
331 modelled as a two-dimensional radial domain, as there is no heterogeneity in the azimuthal
332 direction. The field is segregated into 71x4x10 grids cells. This grid refinement was optimized
333 for less computational time and accurate outputs through a series of experiments that showed
334 no significant effect of the grid refinement up to a magnitude of 10 times on the produced CO₂
335 saturation contours. It is a fact that fine block grids produce smoother contours, however
336 noticeable reduction of execution time was observed by using coarser grids with no momentous
337 influence on the smoothness of CO₂ profiles. This is consistent with studies by Gonzalez-
338 Nicolas et al. (2011) and Hanspal and Das (2012) which indicate that grid refinement has no
339 significant influence on $P_c - S_w$ profiles.

340
341 Supercritical CO₂ is to be injected at pressure and temperature above CO₂ critical conditions
342 into the centre of the computational domain at the lower 30 m at a constant rate of 40 Kg/s for
343 20 years followed by 980 years lockup period as illustrated in Figure 1. This injection rate
344 represents about 25% of a medium size coal-fired power generation plant. Different types of
345 heterogeneities have been considered for the domain and various scenarios of injection
346 process are applied to investigate the effects of permeability, temperature, porosity, and
347 injection pressure on capillary pressure (P_c)-saturation (S_w) relationships at static and dynamic
348 flow conditions which has a significant influence on the fate of CO₂ after the injection process.

349
350 Firstly, we run our simulations on fine and course homogenous domains with porosity of 0.16
351 and 0.25, respectively, to determine the effect of porosity and permeability on $P_c - S_w$
352 relationships. Unlike other works (Mirzaei and Das, 2007; Peszynska and Yi, 2008; Hsu and
353 Hilpert, 2011), dynamic and quasi-static simulations are conducted for comparison purposes.
354 The simulated aquifer and simulation parameters are illustrated in details in Tables 1 and 2,
355 respectively.

356

357 In spite of considering three-dimensional flow in a permeable media, it is noted by Domenico
358 and Schwartz (2000) that under the same hydraulic gradient, horizontal flow is of the order of
359 six orders of magnitude faster than vertical flow. In our simulation the lateral flow dominates,
360 therefore we can refine our computational domain by declaring our reference nodes in the
361 output control card in the input file, vertically at distances from the bottom of the formation at 20,
362 40 and 70 m while horizontally at radial distances of 100, 200, 500, 800, and 1000 m, from the
363 injection well while azimuthally a single reference plane is considered at 45° to measure our
364 simulation variables for two-dimensional scenarios.

365

366 **3. Results and Discussions**

367 To evaluate the behaviour of CO₂ in three-dimensional cylindrical field-scale formation a
368 medium-term of 1000 years of simulations of CO₂ injection into homogeneous and
369 heterogeneous formations were carried out in this work. The injection process continued for 20
370 years followed by 980 years of lockup. Supercritical CO₂ is injected azimuthally at 4 nodes
371 which are uniformly distributed towards the lower 30 m of the domain.

372

373 As stated earlier, this study aims to examine the effects of injection pressure, temperature,
374 heterogeneities (layering), porosity, permeability and injection condition states on the P_c - S_w
375 relationships and the behaviour of the injected CO₂ in terms of its dissolution or mobility. To
376 show how CO₂ behaves over the simulation time period, a series of numerical simulation
377 models displayed in Table 3 were created for different initial and boundary conditions shown in
378 Table 4.

379

380 **3.1. CO₂ Migration**

381 As soon as the injection process starts the supercritical CO₂ displaces the existing brine and
382 migrates away from the injection well as illustrated in Figure 2. For different time levels, the
383 simulated CO₂ spatial distribution profiles are shown for drainage process (Figure 2 (A-D))
384 during the injection period and imbibition process (Figure 2 (E-H)) presenting the post injection
385 period for case 1 simulation conditions (Table 3). CO₂ continues to migrate laterally due to the
386 governing forces, e.g., (i) hydrostatic pressure difference between the injection point and aquifer,
387 and (ii) capillary pressure. Furthermore, as a result of densities difference between the ambient
388 brine and injected supercritical CO₂ buoyancy forces push the latter upwards until it reaches the
389 impervious confining layer (caprock) under which it is trapped or extends more laterally. Figure
390 2 (E-H) demonstrates that when injection process ends the domain is invaded by brine which
391 displaces most of the CO₂ leaving part of it trapped in small pores. This leads to residual
392 trapping of the injected CO₂. Meanwhile a volume fraction of the injected CO₂ are dissolved in

393 the brine after the injection process to produce a rich CO₂ layer that settles permanently at the
394 bottom of the domain.

395

396 **3.2. Effects of Porosity and Permeability**

397 It has been suggested by some researches (e.g., Kumar et al., 2005; Xu et al., 2006; Kopp et
398 al., 2009; Chasset et al., 2011) that an increase in the mean permeability results in greater
399 injectivity and mobility of CO₂ which increases dissolution into the formation brine. To explore
400 this further and, in particular, determine the effects of porosity and permeability on the injectivity
401 of CO₂, two sets of simulations were carried out in this study. The first set in cases 1 and 2
402 explores these effects in fine and coarse homogeneous domains, respectively. The second set
403 (cases 7 and 8) looks at the effects of two heterogeneous porous layers involving fine-coarse-
404 fine and coarse-fine-coarse layers. In all injection cases illustrated in Figure 3, it is observed
405 that the coarse domain produces larger CO₂ plumes during the injection time at middle altitudes
406 of the domain. This means that the higher the permeability the higher the CO₂ saturation is
407 during the drainage process when the hydrostatic forces dominate. Different plumes are
408 produced during the imbibition process when the aquifer brine reverses back to displace the
409 CO₂. Though the CO₂ plume size looks larger for the fine sand domain (Figure 3 (B, C, D)) the
410 actual sequestration of CO₂ was still higher because by then a considerable amount of injected
411 CO₂ had dissolved in the existing brine and most of it would have settled at the top of the
412 domain. This is clearly displayed in the coarse domain contours illustrated in Figure 3 (G,H).

413

414 CO₂ distribution profiles in Figure 4 demonstrates a different behaviour of the injected CO₂ in
415 fine domain where CO₂ residual saturation was never reached though some tendency was
416 noticed at a radial distance of 1000 m after 800 years of simulation compared to the coarse
417 domain where residual CO₂ saturation was reached at about 200 years. In contrast, all CO₂
418 profiles in the coarse domain reach the CO₂ residual saturation levels after 200 years. This is
419 because lower permeability porous media limits both lateral and vertical CO₂ mobility and
420 maintain more contact with the surrounding brine which enhances the solubility trapping to keep
421 the injected gas more securely within the aquifer. In addition, the small size pores in the fine
422 domain play like meniscus tubes which allow CO₂ to break through due to capillary forces to
423 enhance residual trapping. These results are consistent with those obtained by Nordbotten et al.
424 (2005) and Kumar et al. (2005). In the case of heterogeneity these profile show completely
425 different trends which will be discussed in more details in section 3.6.

426

427 **3.3. Effects of Injectivity**

428 Capillarity plays an important role in sequestering CO₂ in geological formations because it
429 enhances the residual trapping (one of the means of sequestration). To investigate the effect of

430 injection pressure on capillarity in the domain, supercritical CO₂ was injected into a coarse sand
431 domain at 36, 34 and 32 MPa (cases 2, 5 and 6), respectively, under dynamic flow conditions.
432 The results presented in Figure 5(A) show no significant influence of the injection pressure on
433 capillary pressure at all saturation values which is most likely due to the employed values of
434 injection pressure being very close to the hydrostatic pressure in the aquifer, in fact in case 6
435 the same value of 32 MPa was used in addition to the high permeability of the domain which
436 offers easier migration of CO₂ laterally and vertically. These results are qualitatively consistent
437 with the experimental results achieved by Plug and Brunning (2007).

438
439 The influence of injection pressure on CO₂ saturation profiles is displayed in Figure 5(B), which
440 shows a steep increase in CO₂ saturation at 70 m altitude for all injection pressures as a result
441 of gravity forces which cause most of the injected CO₂ to migrate up towards the top of the
442 aquifer. This increase reaches the highest value at 200 years when the trends sharply steep
443 down till they reach CO₂ residual saturation. The results indicate that the higher the injection
444 pressure, the larger the amount of CO₂ accumulated at the top of the aquifer at times between
445 200 – 500 years of simulation. This amount of CO₂ is unlikely subject to any permanent trapping
446 in short-term periods of simulation because it is not affected by the imbibition process.

447 448 **3.4. Temperature Effects on CO₂ Distribution**

449 In this research work we inspected the effects of temperature on CO₂ distribution during
450 drainage and imbibition processes. CO₂ saturation contours in Figure 6 demonstrate smaller
451 plumes of supercritical CO₂ after 20 years of injection (i.e. end of drainage process) and 200
452 years (during imbibition process) at a domain temperature of 80°C (case 5) compared to those
453 for 58°C (case 3) under the same injection pressure conditions. This is because increasing the
454 temperature decreases the density and viscosity of the injected CO₂ and consequently
455 increases buoyancy and gravity forces that contribute in spreading CO₂ laterally and vertically.
456 The effect of temperature on CO₂ dissolution in the hosted brine is illustrated in Figure 7. It
457 shows that after about 200 years of simulation higher temperature results in more CO₂
458 dissolved because it decreases CO₂ density which migrates upwards to get in contact with
459 more fresh brine which enhances the solubility trapping mechanism. Moreover it is observed in
460 Figure 8 that capillary pressure increases proportionally with temperature and this change is
461 more prominent between saturation values of 0.55 - 0.7. The results in Figure 8 demonstrate
462 that about 50% increase in capillary pressure is obtained when the temperature is increased
463 from 58°C to 80°C. This increase in capillary pressure permits more CO₂ flow into the small
464 pores and trap there as a residual solute, which is referred to as residual trapping of CO₂.

465

466 **3.5. Dynamic Capillary Pressure Effects**

467 Our investigations explain that in addition to the saturation, capillary pressure is strongly
468 influenced by the flow conditions in the system. Several simulation tests were carried out to
469 compare the CO₂ saturation change in homogeneous and heterogeneous computational
470 domains under quasi-static and dynamic conditions. The results are illustrated in Figure 9 which
471 demonstrates that higher CO₂ saturation is obtained under quasi-static flow conditions at any
472 radial distance from the injection well at all time levels. The elapsed time to attain static
473 conditions allows more CO₂ into small pores by capillary forces and it may increase convective
474 mixing between the two fluids which enhances solubility trapping of the injected CO₂.
475 Additionally, it is noticed from Figure 10 that at saturation values above 0.55 higher capillary
476 pressures are generated in fine-grained domain which is consistent to the theories which relate
477 capillary pressure directly to the pore size.

478
479 To determine the dynamic or damping coefficient (τ), which indicates the extent of dynamic
480 capillary pressure effect, two numerical simulations (cases 2 and 10) were run under dynamic
481 and quasi-static condition, respectively. All calculation results are displayed in Figure 11. From
482 equation (16) we calculated the dynamic coefficient for each average value of the aqueous
483 saturation calculated by equation (14), the corresponding average values of dynamic and static
484 capillary pressures ($P_{c,dyn}$, $P_{c,stat}$), and the calculated values of the time derivative of saturation
485 ($\partial S/\partial t$). Figure 11 shows that the value of dynamic coefficient decreases when the rate of
486 change of aqueous saturation increases and this decline is very sharp at low saturation values
487 when the rate of change in aqueous saturation is slow (i.e. lower values of $\partial S/\partial t$). The attained
488 relationship between the dynamic coefficient and aqueous saturation can be clarified by the
489 longer time required to attain the residual saturation at higher values of dynamic coefficient.

490

491 **3.6. Effects of Heterogeneity**

492 Heterogeneity is closely related to the disparity in permeability which strongly rules the CO₂
493 transport through different parts of the domain. To investigate this influence we compare four
494 study cases (1, 2, 7 and 8) in terms of CO₂ saturation distribution in homogenous and
495 heterogeneous domains. As expected, and in agreement with previous studies (e.g., Ataie-
496 Ashtiani et al., 2001; Das et al., 2006), heterogeneity has shown an important influence on the
497 characteristics of two-phase flow. It is shown in Figure 12 that all trends behave similarly at 200
498 m and 40 m horizontal and vertical distances, respectively, as they display an increase in the
499 integrated aqueous CO₂ upon injection stops and tend to plateau after about 200 years except
500 in case 8 (coarse sand embedded in fine) which starts to drop after 50 years of simulation. This
501 behaviour is related to the injection section into which the supercritical CO₂ was injected (lower
502 30 m) which for this case is a fine layer bounded by a coarse one above. This scenario

503 encourages vertical migration of CO₂ due to the lower entry pressure of the upper strata which
504 consequently reduces both residual and solubility trapping by eliminating the contact time
505 between CO₂ and existing brine. However, larger amount of CO₂ was dissolved in the
506 homogenous coarse domain due to the high permeability which increases CO₂ movement in
507 both directions maintaining more contact with fresh brine in which it dissolves. Figure 13
508 presents the total integrated amount of CO₂ in aqueous and gas phases at the same grid block
509 described above. It is apparent that all curves decline and tend to plateau soon after the
510 injection stops except case 8 which shows sharp increase in total integrated CO₂ till about 50
511 years of simulation and continuously increases until the end of 1000 years of simulation. This
512 can be explained by the fact that the injected CO₂ favourably move through large size pores
513 which increases the hydrodynamic trapping as a result of pressure difference forces and CO₂
514 concentration. This is combined with the solubility trapping due to the convective mixing of CO₂
515 and the surrounding brine. Residual trapping is larger in the surrounded fine layer which slows
516 down the migration of the injected CO₂ providing more chance to enter the small pores.

517
518 CO₂ spatial spread is demonstrated in three-dimensional cylindrical contours in Figure 14,
519 which demonstrates how the injected supercritical CO₂ spreads through different
520 heterogeneous domains (cases 7 and 8) at different time steps. During a drainage process
521 period of 20 years for case 7 (coarse sand embedded in fine), higher CO₂ saturation values
522 were obtained. This increase is a result of the pressure difference forces that control the lateral
523 migration of CO₂ and vertically due to the buoyancy forces that transfer the supercritical fluid
524 from the low permeability layer up to the higher permeability one as evidently shown in Figures
525 14(A-D).

526
527 In contrast during imbibition process (post injection) which is presented in case 8, higher
528 concentrations of CO₂ were achieved as shown in Figures 14(E,H) because in this case
529 advective, diffusive and gravity forces all contribute in trapping the injected CO₂ in addition to
530 the reversed-back movement of the brine behind the CO₂ plume to displace the CO₂ again
531 leaving traces of it as residual contaminants in small-sized pores which is referred to as residual
532 trapping. Moreover the low permeability layer retains more contact time between the two fluids
533 which enhances the solubility trapping mechanism. The results of this work are consistent with
534 those obtained by Nordbotten et al. (2005) as they found out that the buoyancy forces places
535 the highest mobility layer of injected CO₂ at the top of the domain considering that in their case
536 the CO₂ was injected along the whole altitude of the domain.

537

538 **4. Conclusion**

539 A series of numerical simulations have been conducted in this work to identify the possible
540 implications of a number of important parameters on the capillary pressure – saturation
541 relationship for supercritical CO₂ in deep saline aquifer. From the results of this work it is
542 obvious that the higher the injection pressure, the higher the capillary forces are; however, the
543 maximum sustainable pressure has to be taken into consideration to avoid any geochemical
544 fracture to the formation rock. The value of the dynamic coefficient (τ) increases as the rate of
545 change of aqueous saturation ($\partial S/\partial t$) declines because more time is required for the residual
546 saturation to be attained. It has been found that capillary forces are higher in fine-grained
547 domains and they enhance storage capacity of the site by amplifying the residual trapping
548 mechanism of CO₂ during the imbibitions process. Solubility trapping is more efficient in fine
549 domains because they maintain more contact between the fluid phases which leads to more
550 CO₂ dissolved in aquifer brine. Warm aquifers are more competent in CO₂ sequestration
551 because higher temperatures increase the capillary pressure and consequently enhance
552 residual trapping of CO₂. Fine sand embedded in coarse pattern of heterogeneity is found to be
553 more effective method over long periods of storage procedure however more research is
554 required to clarify how the field distribution of heterogeneity and injection scenarios of CO₂
555 affect the efficiency of the sequestration.

556

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560

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Table 1. Selected aquifer parameters for simulation

Parameter	Value/Function		Reference
Diameter (m)	5000		-
Thickness (m)	100		-
Depth (m)	2900		-
Grid (nodes)	71 x 4 x 10		-
Porosity	Sandstone 0.25	(Wechselfolgen) 0.16	May et al. (2004)
Horizontal Permeability (m ²)	(5.625e-13)	(0.5428e-13)	May et al. (2004)
Vertical Permeability (m ²)	(1.688e-13)	(11.15e-16)	May et al. (2004)
rock density (kg/m ³)	2430	2470	May et al. (2004)
specific storativity (1/m)	9.2e-4		May et al. (2004)
Surface temperature (°C)	8		May et al. (2004)
Reservoir temperature (°C)	58		Beni et al. (2012)
Temperature gradient (K/m)	0.035		Reinicki, (1968)
Reservoir pressure (MPa)	32		Beni et al. (2012)
Pressure gradient (KPa/m)	10.5		Beutler, (1975)

Table 2. Important parameters and initial conditions

Parameter	Value	Reference
Irreducible saturations water, S_{ir}	0.1	Beni et al. (2012)
CO ₂ , S_{gr}	0.05	Beni et al. (2012)
Brooks/Corey Exponent, λ	0.457	Beni et al. (2012)
Strength coefficient, P_0	19,610 Pa	Beni et al. (2012)
Pore compressibility, k	$1 \times 10^{-9} \text{ Pa}^{-1}$	Beni et al. (2012)
Pore expansivity, β	$1 \times 10^{-6} \text{ K}^{-1}$	Beni et al. (2012)
Injection Pressure	36 MPa	-
Temperature	58 °C	Beni et al. (2012)
Salinity	0.2	Beni et al. (2012)
Pressure gradient	10.5 MPa/Km	May et al. (2004)
Salinity gradient	80 g/L.Km	May et al. (2004)
CO ₂ injection rate	40 Kg/s	-
Injection time	20 Yrs	-
Simulation time	1000 Yrs	-

Table 4. Initial and boundary conditions

Case No.	Domain Type/Cond.	Horizontal Permeability (m ²)	Domain Temp. (°C)	CO ₂ Injection Pressure (MPa)
<i>Dynamic</i>				
1	Homogenous Fine Sand	0.5428e-13	58	36
2	Homogenous Coarse Sand	5.625e-13	58	36
3	Homogenous Coarse Sand	5.625e-13	70	36
4	Homogenous Coarse Sand	5.625e-13	80	36
5	Homogenous Coarse Sand	5.625e-13	58	34
6	Homogenous Coarse Sand	5.625e-13	58	32
7	Heterogeneous Coarse in Fine Sand	0.5428e-13 - 5.625e-13 - 0.5428e-13	58	36
8	Heterogeneous Coarse in Fine Sand	5.625e-13 - 0.5428e-13 - 5.625e- 13	58	36
<i>Quasi-Static</i>				
9	Homogenous Fine Sand	0.5428e-13	58	36
10	Homogenous Fine Sand	0.5428e-13	58	36
11	Homogenous Fine Sand	0.5428e-13	58	32
Porosity; Fine Sand = 0.16, Coarse Sand = 0.25				
Vertical Permeability; Fine Sand = 0.01115e-13, Coarse Sand = 1.6876e-13				
Hydrostatic Pressure=32 MPa, Pressure Gradient= 10.5 MPa/Km, Salinity=0.2				

Source: illustrated in Tables 1 and 2

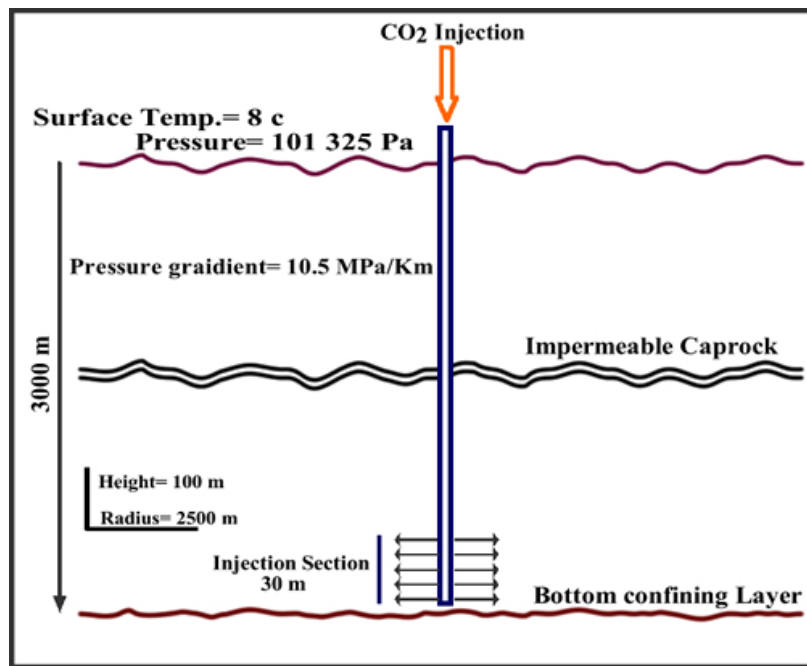


Figure 1. A schematic diagram of geological CO₂ sequestration process in a deep saline aquifer (DSA).

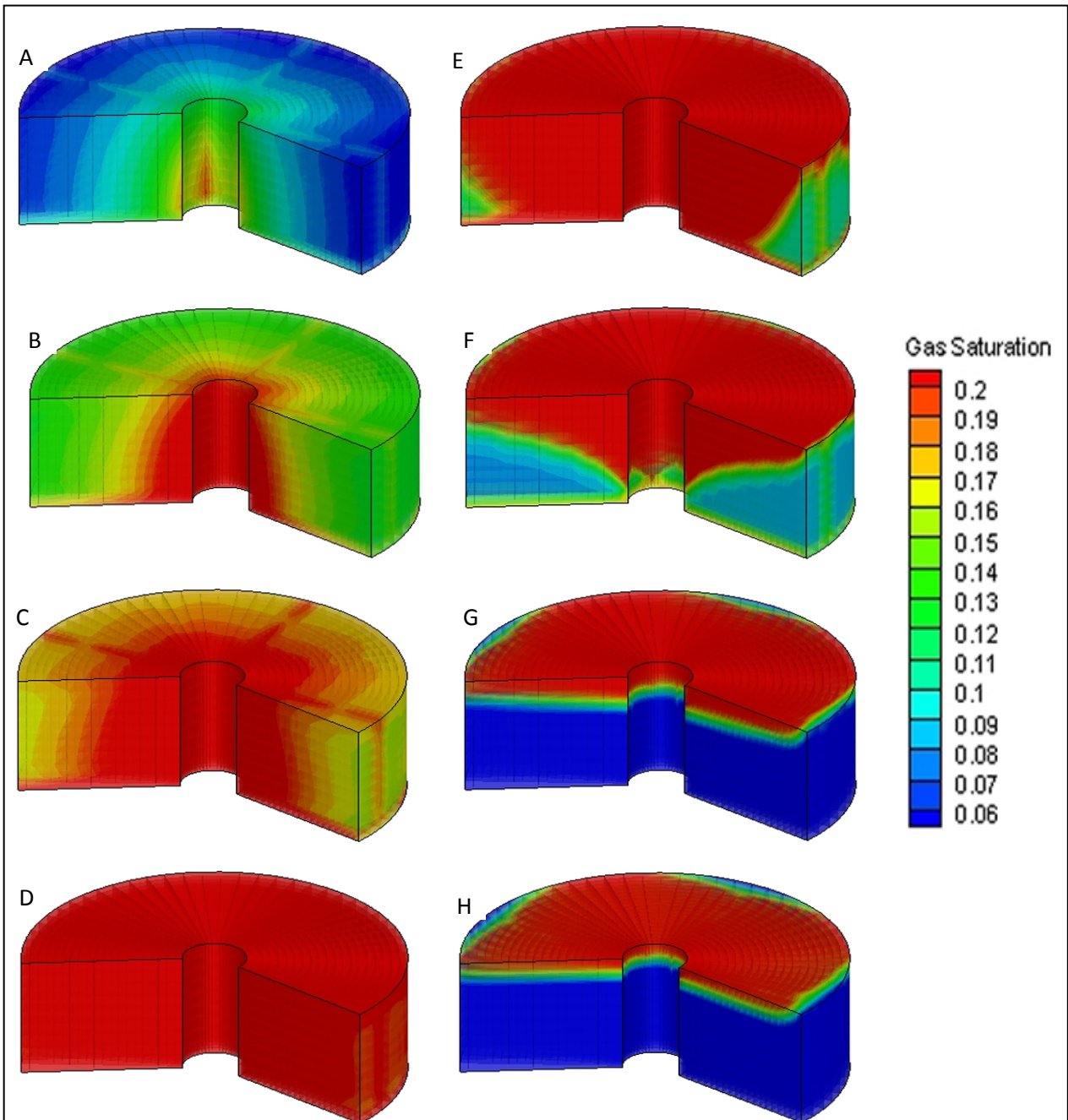


Figure 2. Evolution of CO₂ plume in low permeability homogenous aquifer (case1):

(A-D) CO₂ distribution after 1, 5, 10 and 20 years of simulation during injection process (drainage).
 (E-H) CO₂ distribution after 50, 200, 500 and 1000 years post injection process (imbibition).

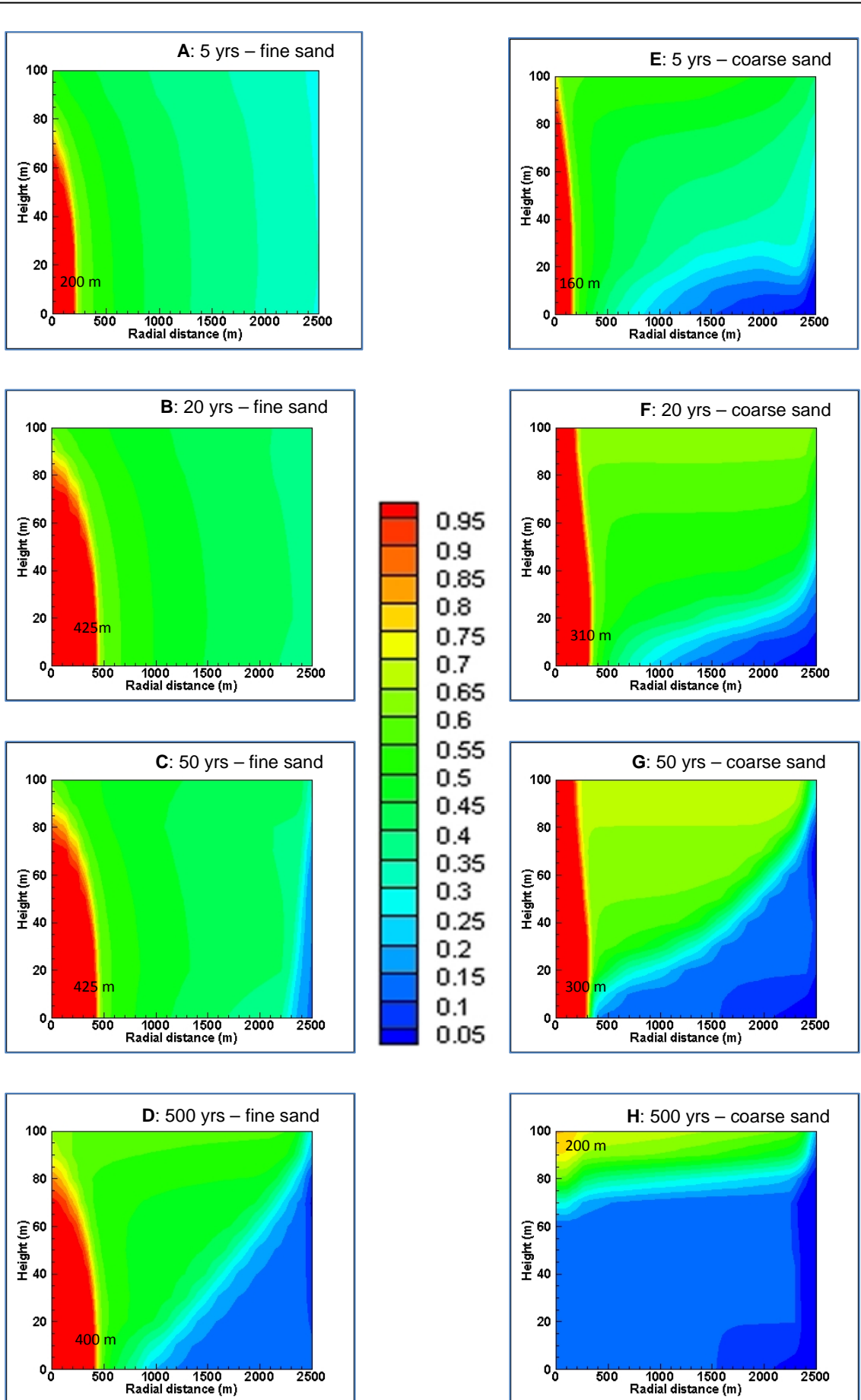


Figure 3. CO₂ plume evolution in low and high permeability homogeneous aquifers at different time scales; (A-D) fine sand domain (Case 1), (E-H) coarse sand domain (Case 2).

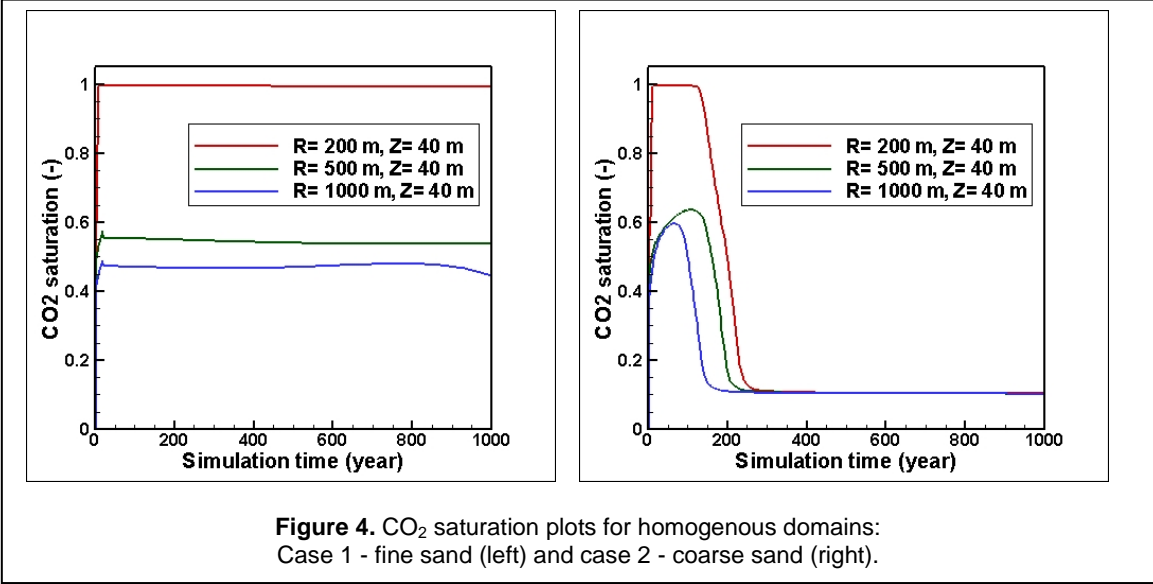


Figure 4. CO₂ saturation plots for homogenous domains:
Case 1 - fine sand (left) and case 2 - coarse sand (right).

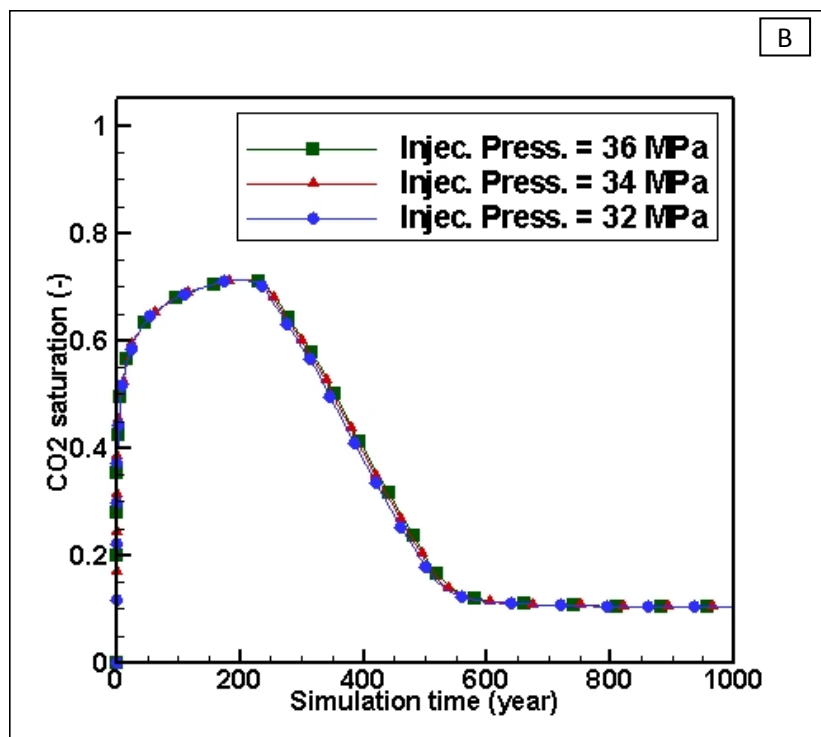
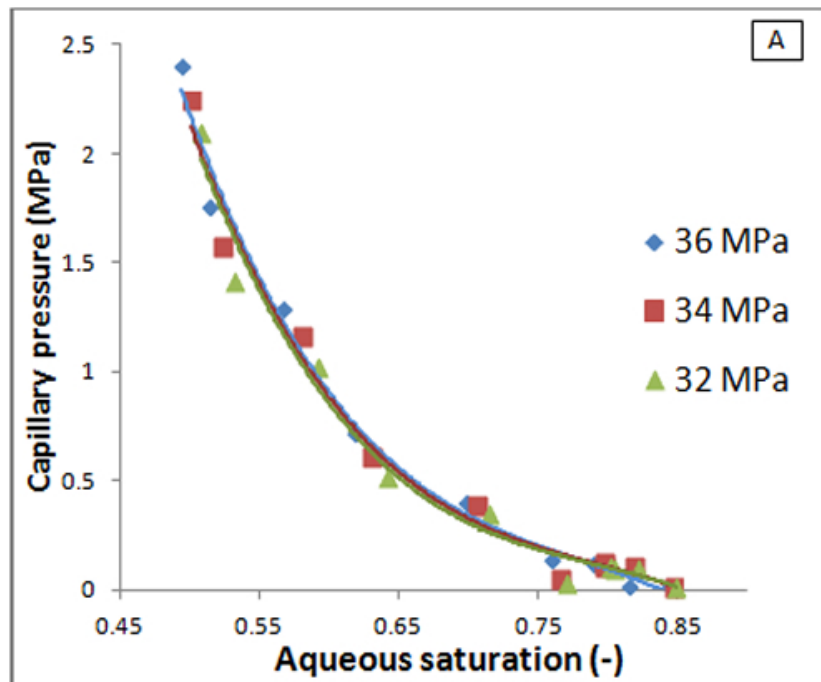
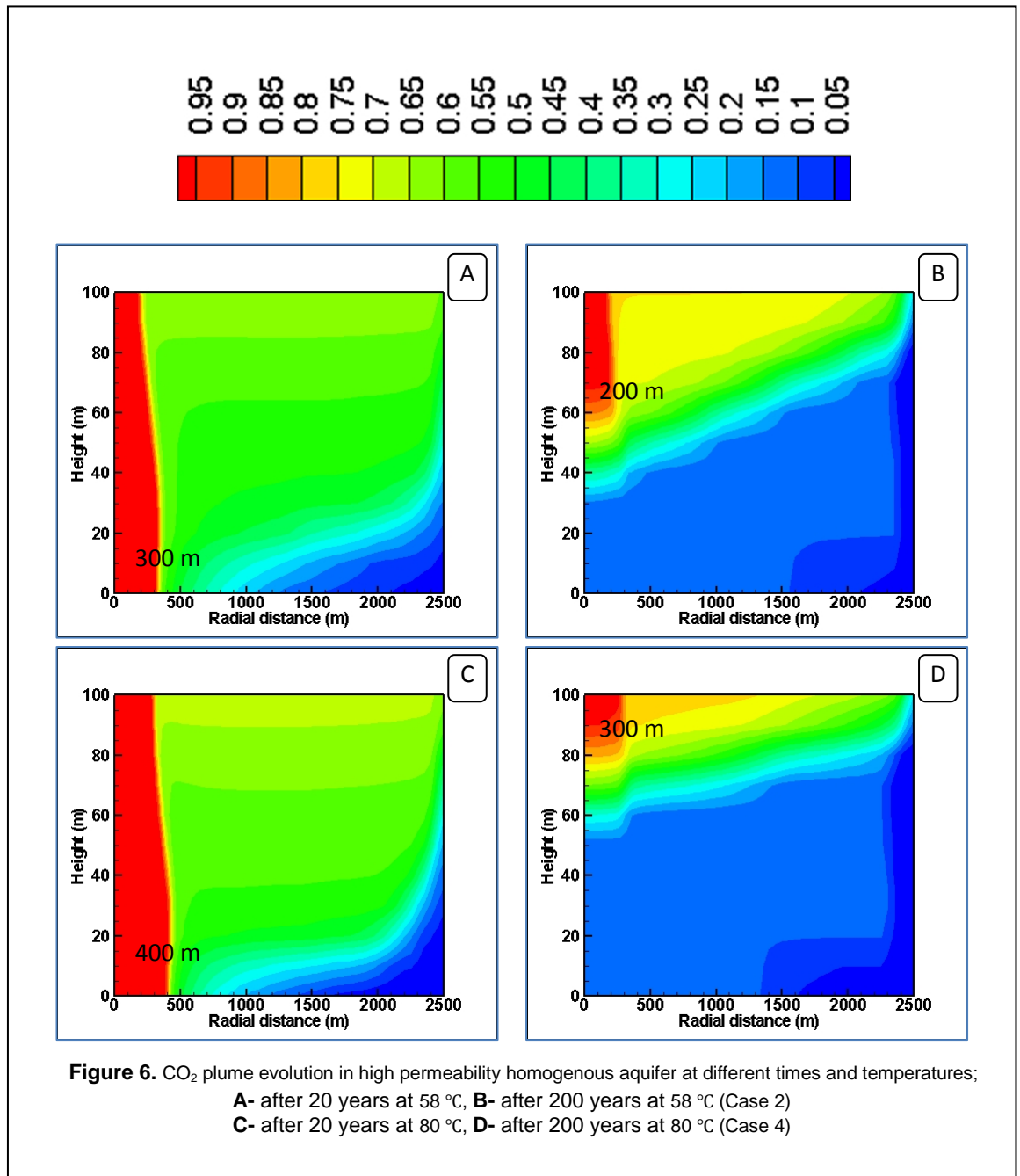


Figure 5. A- Capillary pressure vs. aqueous saturation at different injection pressures.
 B- CO₂ saturation curves at R=500 m and Z=70 m in a coarse homogeneous domain.



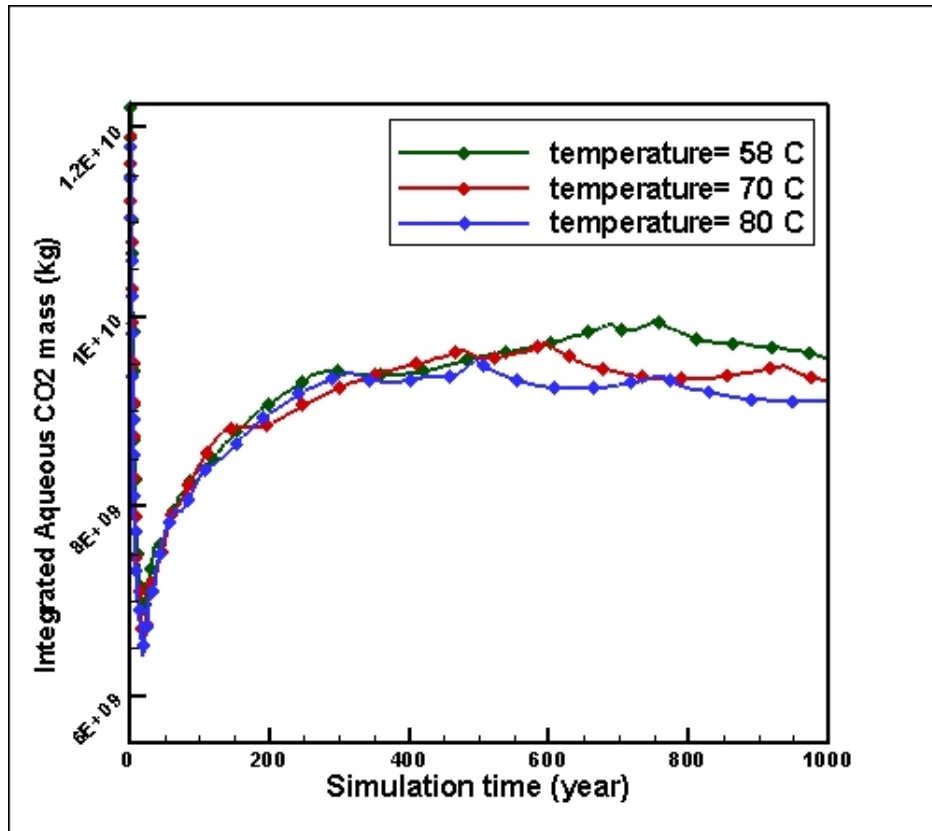


Figure 7. Temperature effects on dissolved CO₂ mass (cases 2, 3 and 4).

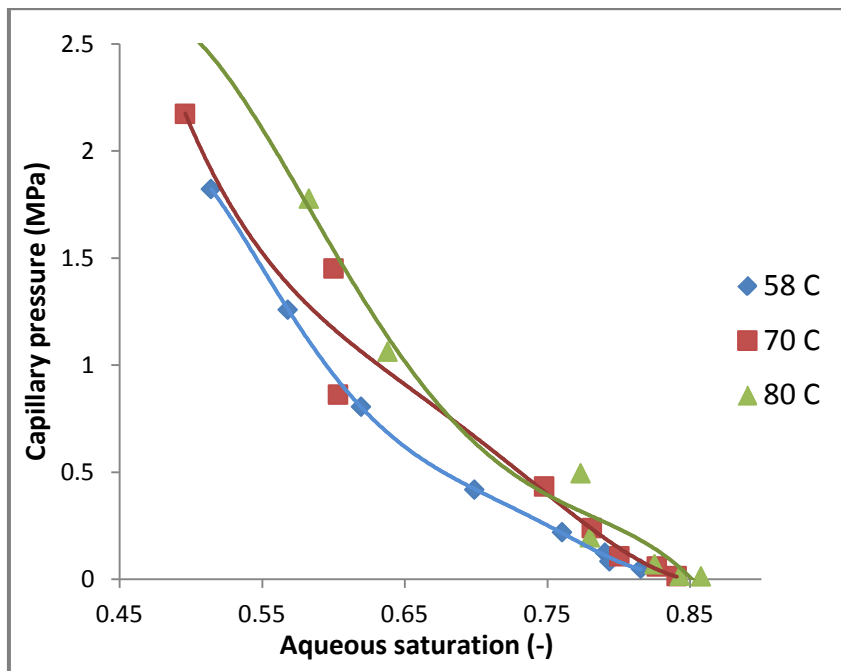


Figure 8. Capillary pressure vs. aqueous saturation at different temperatures (cases 2, 3 and 4).

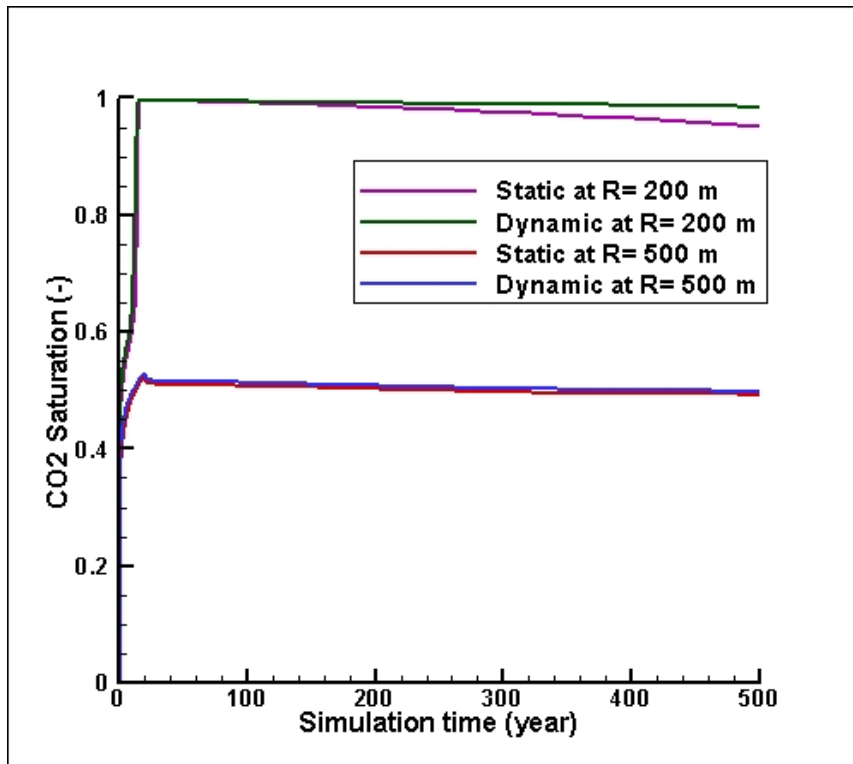


Figure 9. CO₂ saturation (volume fraction) curves for homogenous fine domain under static and dynamic conditions at altitude of 40 m (cases 1 and 9).

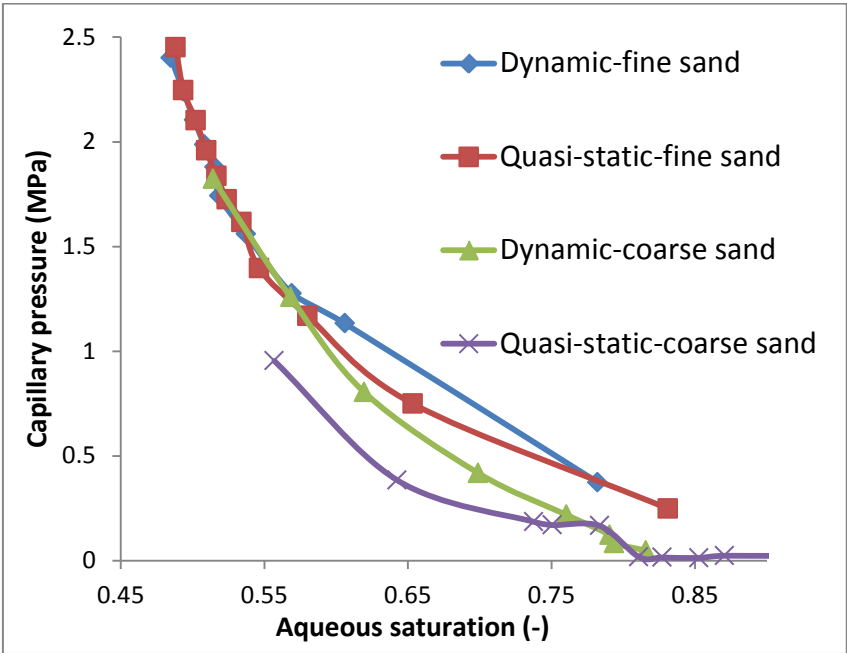


Figure 10. Dynamic and quasi-static capillary pressure-saturation curves for homogenous domains.

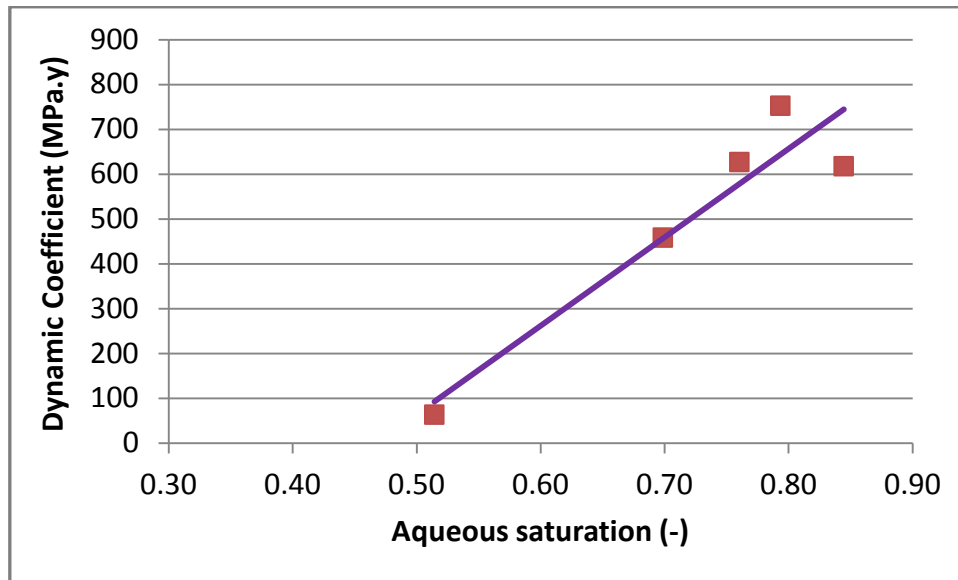


Figure 11. Dynamic Coefficient change with aqueous saturation in a homogeneous coarse sand domain (cases 2 and 10).

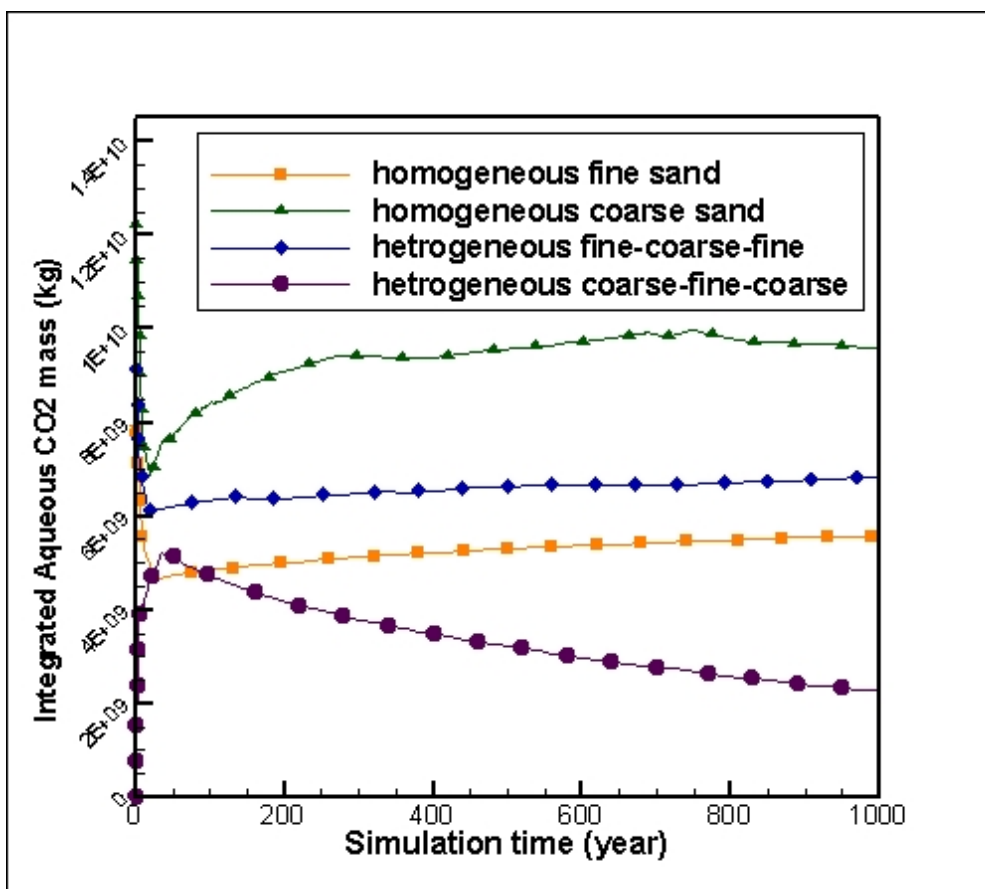


Figure 12. Integrated aqueous change for homogeneous and heterogeneous domains (cases 1, 2, 7 and 8) at radial distance of 200 m and altitude of 40 m

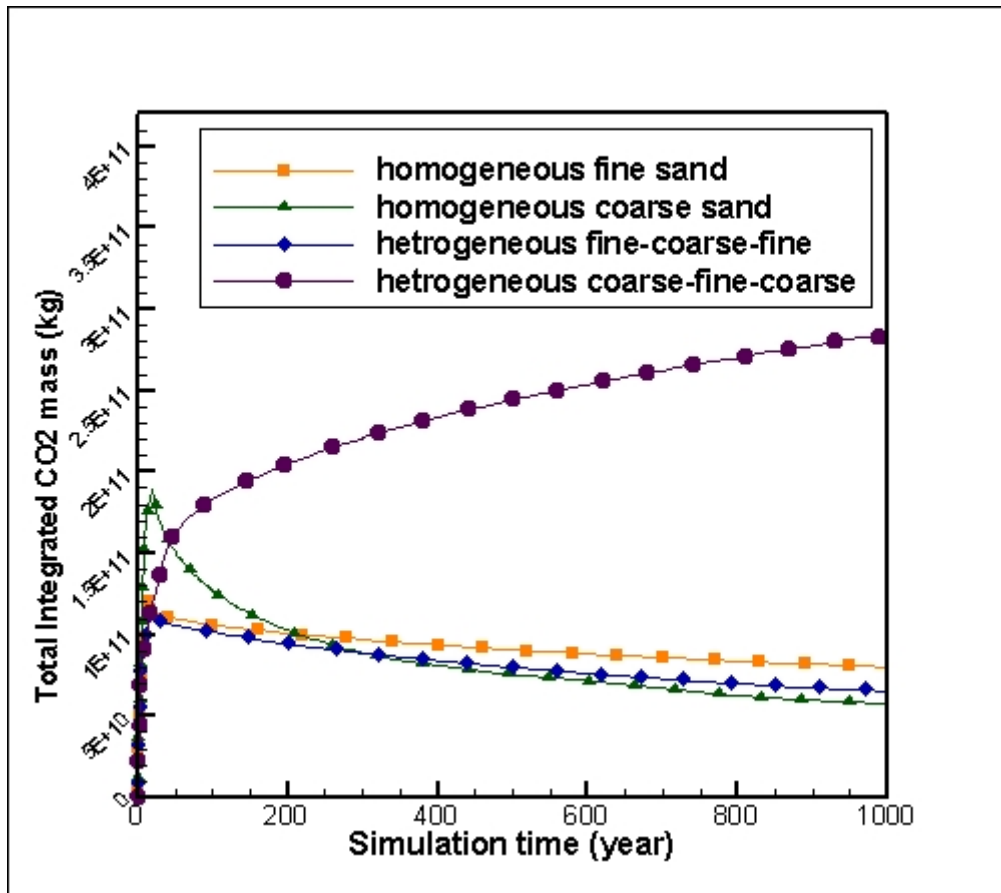


Figure 13. Total Integrated CO₂ profiles for homogeneous and heterogeneous domains (cases 1, 2, 7 and 8) at radial distance of 200 m and altitude of 40 m.

